

# Crystal structure and Mössbauer effect studies of $\text{Gd}(\text{Mn}_{1-x}\text{Fe}_x)_2$ and $\text{Gd}(\text{Fe}_{1-x}\text{Co}_x)_2$ intermetallics

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**Abstract.** X-ray crystal structure (300 K) and  $^{57}\text{Fe}$  Mössbauer effect (77 K) studies were performed for the  $\text{Gd}(\text{Mn}_{1-x}\text{Fe}_x)_2$  and  $\text{Gd}(\text{Fe}_{1-x}\text{Co}_x)_2$  compounds. The crystal lattice parameter decreases vs. an average number of  $3d$  electrons. The magnetic hyperfine field treated as a function of the average number of  $3d$  electrons forms a Slater-Pauling type dependence. This dependence is related to the  $3d$ -electron band structure within the frame of the rigid band model.

**Key words:** intermetallics • crystal structure •  $^{57}\text{Fe}$  Mössbauer effect • hyperfine interactions • Slater-Pauling dependence • band structure

## Introduction


Fundamental interest and practical applications are the reason behind the numerous studies of the heavy rare earth (R) – transition metal (M) compounds [1, 2, 6, 7]. It was previously found that the hyperfine interactions and the magnetic properties of the R-M materials are mainly governed by the  $3d$  electrons of the M sublattice [3, 5].

Thus, it was interesting to study the significance of the Mn/Fe and Fe/Co substitution in the  $\text{Gd}(\text{Mn}_{1-x}\text{Fe}_x)_2$  and  $\text{Gd}(\text{Fe}_{1-x}\text{Co}_x)_2$  intermetallics. For this purpose, crystal structure and  $^{57}\text{Fe}$  Mössbauer effect studies were performed. The obtained data are qualitatively discussed within the frame of the rigid band model [8].

## Experiment

The intermetallics  $\text{Gd}(\text{Mn}_{1-x}\text{Fe}_x)_2$  and  $\text{Gd}(\text{Fe}_{1-x}\text{Co}_x)_2$  ( $0 \leq x \leq 1$ ) were prepared by arc melting, in a high purity argon atmosphere from the appropriate amounts of Gd (99.9% purity), Mn, Fe and Co (all 99.99% purity).

From the X-ray patterns, the cubic,  $Fd\bar{3}m$ ,  $\text{MgCu}_2$ -type (C15) Laves phases were deduced for all the compounds. Since the atomic radius is reduced across the transition metal series ( $r_{\text{Mn}} = 1.79 \text{ \AA}$ ,  $r_{\text{Fe}} = 1.72 \text{ \AA}$  and  $r_{\text{Co}} = 1.67 \text{ \AA}$ ) thus the lattice parameter  $a(n)$  (Fig. 1) is reduced vs. the average number  $n$  of  $3d$  electrons calculated per transition metal atom. A typical deviation for R-M intermetallics from Vegard's rule is observed also for studied compounds. Line in Fig. 1 following experimental points is described by a numerical formula  $a(n) = (-0.109n^3 + 2.062n^2 - 13.123n + 35.440) \text{ \AA}$ .

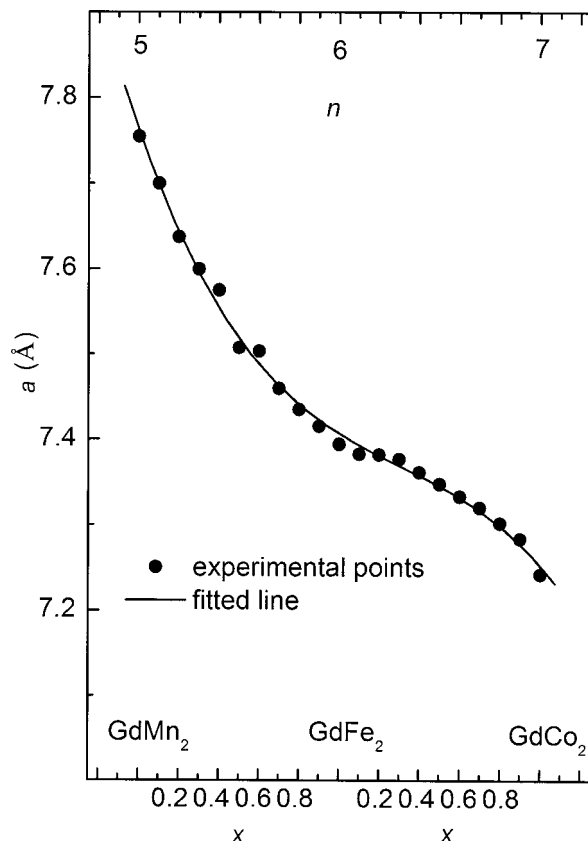
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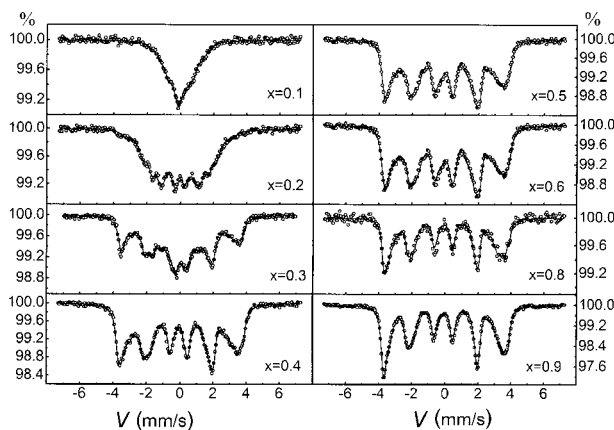
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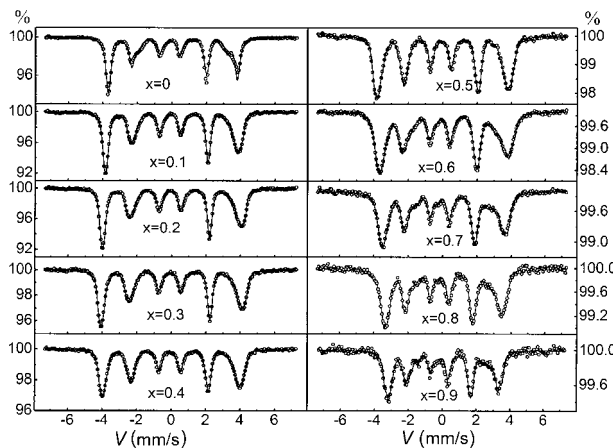


**Fig. 1.** The crystal lattice parameters of the  $\text{Gd}(\text{Mn}_{1-x}\text{Fe}_x)_2$  and  $\text{Gd}(\text{Fe}_{1-x}\text{Co}_x)_2$  compounds (300 K) vs.  $n$  the average number of  $3d$  electrons.

Mössbauer spectra were collected using a standard transmission technique with a Co/Rh source. The resulting  $^{57}\text{Fe}$  Mössbauer effect spectra (at 77 K) for the  $\text{Gd}(\text{Mn}_{1-x}\text{Fe}_x)_2$  and  $\text{Gd}(\text{Fe}_{1-x}\text{Co}_x)_2$  series are presented in Figs. 2 and 3, respectively. Fitting procedure was performed considering both the  $\langle 111 \rangle$  easy axis of magnetization (two subspectra with the ratio 0.25/0.75) and the assumed random distribution of transition metal atoms (Bernoulli distribution) by an analogous method as presented previously elsewhere [4]. Figure 4 contains the determined average hyperfine

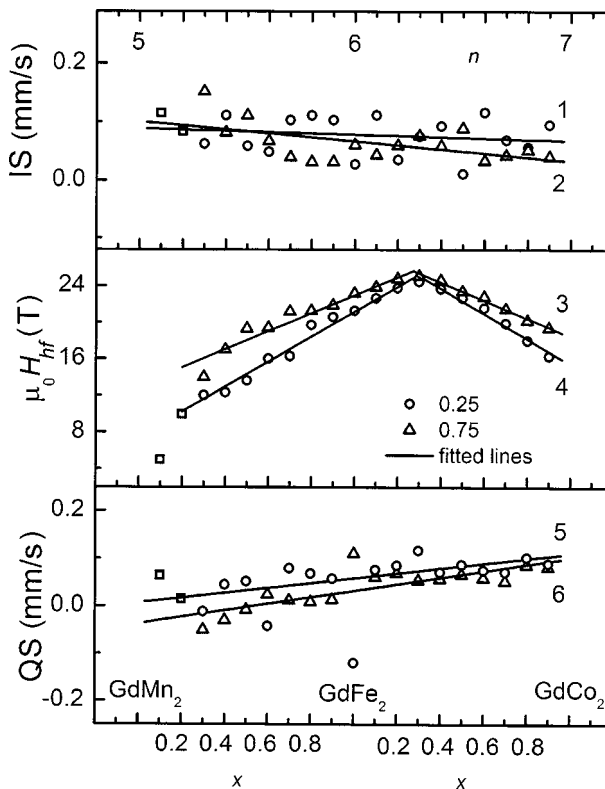


**Fig. 2.**  $^{57}\text{Fe}$  Mössbauer effect spectra of the  $\text{Gd}(\text{Mn}_{1-x}\text{Fe}_x)_2$  and  $\text{Gd}(\text{Fe}_{1-x}\text{Co}_x)_2$  intermetallics (77 K). Experimental points and fitted lines are presented.



**Fig. 3.**  $^{57}\text{Fe}$  Mössbauer effect spectra of the  $\text{Gd}(\text{Fe}_{1-x}\text{Co}_x)_2$  intermetallics (77 K). Experimental points and fitted lines are presented.

interaction parameters, i.e. the isomer shift IS (related to pure iron metal at 300 K), the magnetic hyperfine field  $\mu_0 H_{hf}$  ( $\mu_0$  is the magnetic permeability) and the quadrupole interaction parameter QS (as defined for instance in [6]). The dependencies (1, 4, 5) and (2, 3, 6) correspond to the 0.25 and the 0.75 subspectra, respectively. As a result of considerable experimental errors, as compared to small IS and QS values lines 1, 2 and 5, 6 are weakly resolved. Nevertheless, curves 3 and 4 are resolved satisfactorily. Both these curves follow a Slater-Pauling type dependence.



**Fig. 4.** Average hyperfine interaction parameters of the  $\text{Gd}(\text{Mn}_{1-x}\text{Fe}_x)_2$  and  $\text{Gd}(\text{Fe}_{1-x}\text{Co}_x)_2$  series (77 K): (1, 4, 5) and (2, 3, 6) the hyperfine interaction parameters for the 0.25 and the 0.75 subspectra, respectively.

## Discussion

The substitution of one transition metal by another transition metal strongly influences the magnetic hyperfine fields. As a result of the Mn/Fe and Fe/Co substitution, the field  $\mu_0 H_{hf}(n)$  creates the Slater-Pauling type curve. At first, the strong ferromagnetism-type behavior of the M-sublattice is observed. The magnetic hyperfine field in the  $\text{Gd}(\text{Mn}_{1-x}\text{Fe}_x)_2$  series grows up with  $x$ . This growth occurs for the  $\text{Gd}(\text{Fe}_{1-x}\text{Co}_x)_2$  series too, and the maximum value of the field is approached for the  $\text{Gd}(\text{Fe}_{0.7}\text{Co}_{0.3})_2$  compound (at  $x = 0.3$ ,  $n = 6.3$ ). At this Co-content, the filling up of the majority 3d-subband by 3d-electrons is terminated. For the higher Co-substitutions, the weak ferromagnetism-type behavior of the M-sublattice appears. The filling up of the minority 3d-subband still proceeds, and the observed field decreases gradually with  $n$ . The  $\mu_0 H_{hf}$  field strongly depends on the average number  $n$  of 3d electrons calculated per transition metal atom.

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